

**O‘ZBEKISTON RESPUBLIKASI OLIY TA’LIM, FAN VA INNOVATSIYALAR  
VAZIRLIGI**

**MIRZO ULUG‘BEK NOMIDAGI O‘ZBEKISTON MILLIY UNIVERSITETINING  
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**RAQAMLI JAMIYATDA FAN VA TA’LIMNING INTEGRATSIYALALASHUVI:  
MATEMATIKA, IQTISOD, PEDAGOGIKA VA PSIXOLOGIYANING YANGI  
YONDASHUVLARI**  
*mavzusidagi Respublika ilmiy-texnik anjuman materiallari to‘plami*  
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**Raqamli jamiyatda fan va ta’limning integratsiyalalashuvi: matematika, iqtisod, pedagogika va psixologiyaning yangi yondashuvlari.** Respublika ilmiy-texnik anjuman materiallari to’plami – Jizzax: O‘zMU Jizzax filiali Amaliy matematika kafedrası, 2026-yil 20-21-aprel 505-bet.

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Mas’ul muharrirlar: DSc.prof. Turakulov O.X., t.f.n., dots. Baboyev A.M.

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## ON POSITIVE EIGENVALUES OF THE DISCRETE SCHRÖDINGER OPERATOR WITH A NON-LOCAL POTENTIAL

Lakaev Shukhrat, National university of Uzbekistan, Phd, Docent

Tel. +998977038516, [shlakaev85@gmail.com](mailto:shlakaev85@gmail.com)

**Abstract.** On the  $d$  – dimensional lattice  $Z^d$ ,  $d \geq 1$  the discrete Schrödinger operator  $H_{\lambda\mu}$  with a non-local potential constructed via the Dirac delta function and shift operator considered. The dependence of positive eigenvalues existence of the operator on the parameters is explicitly derived.

**The discrete Schrödinger operator in the position representation.** The standard discrete Laplacian  $\Delta$  on the lattice  $Z^d$ ,  $d \geq 1$ , is defined with the following self-adjoint (bounded) multidimensional Toeplitz-type operator on the Hilbert space  $\ell^2(Z^d)$  ([3])

$$-\Delta = \frac{1}{2} \sum_{s \in Z^d, |s|=1} (T(0) - T(s)),$$

where  $T(y)$ ,  $y \in Z^d$  is the shift operator

$$(T(y)f)(x) = f(x + y), \quad f \in \ell^2(Z^d).$$

Let  $V_0$  be a multiplication operator in  $\ell^2(Z^d)$  by the Kronecker delta function  $\delta[\cdot, 0]$

$$V_0 f(x) = \delta[x, 0] f(x).$$

Then, for a given point  $x_0 \in Z^d$ , we define the potential [3] as

$$\tilde{V}_{x_0} = \mu V_0 + \lambda (V_0 T(x_0) + T^*(x_0) V_0).$$

The discrete Schrödinger operator  $H_{\lambda\mu}$  acting in  $\ell^2(Z^d)$ , in the position representation, is defined as a bounded self-adjoint perturbation of  $-\Delta$  and is of the form  $H_{\lambda\mu} = -\Delta - V$ .

**The discrete Schrödinger operator in momentum representation.** The one-particle Hamiltonian  $H_{\lambda\mu}$  in the momentum representation is given as

$$H_{\lambda\mu} = H_0 - V,$$

where  $H_0$  is defined as

$$H_0 = F^*(-\Delta)F,$$

where  $F$  stands for the standard Fourier transform  $F : L^2(T^d) \rightarrow \ell^2(Z^d)$  with the inverse  $F^* : \ell^2(Z^d) \rightarrow L^2(T^d)$ . Explicitly, the non-perturbed operator  $H_0$  acts on  $L^2(T^d)$  as a multiplication operator by the function  $e(p)$

$$(H_0 f)(p) = e(p)f(p), \quad f \in L^2(T^d),$$

where

$$e(p) = \sum_{j=1}^d (1 - \cos p_j), \quad p \in T^d.$$

In the physical literature, the function  $e(p)$  being a real valued-function on  $T^d$ , is called the *dispersion relation* of the Laplace operator.

The perturbation  $V$  is the two-dimensional integral operator

$$(Vf)(p) = (V_{x_0} f)(p) = \frac{1}{(2\pi)^d} \int_{T^d} \left( \mu + \lambda(e^{i(x_0,p)} + e^{-i(x_0,s)}) \right) f(s) ds, \quad f \in L^2(T^d).$$

**The essential spectrum.** The perturbation  $V$  of  $H_0$  is a two dimensional operator, therefore in accordance with the Weyl theorem on the stability of the essential spectrum the equality  $\sigma_{\text{ess}}(H_{\lambda\mu}) = \sigma_{\text{ess}}(H_0)$  holds, and the essential spectrum  $\sigma_{\text{ess}}(H_{\lambda\mu})$  of the operator  $H_{\lambda\mu}$  fills in the following interval on the real axis

$$\sigma_{\text{ess}}(H_{\lambda\mu}) = [e_{\min}, e_{\max}],$$

where

$$e_{\min} = 0, \quad e_{\max} = 2d.$$

The Fredholm determinant of  $H_{\lambda\mu}$ . Set

$$a(z) = \frac{1}{(2\pi)^d} \int_{\mathbb{T}^d} \frac{1}{e(t) - z} dt, \quad b(z) = \frac{1}{(2\pi)^d} \int_{\mathbb{T}^d} \frac{e^{i(x_0, t)}}{e(t) - z} dt, \quad z \in \mathbb{R} \setminus [e_{\min}, e_{\max}].$$

For any  $\lambda, \mu \in \mathbb{R}$ , the Fredholm determinant associated to the operator  $H_{\lambda\mu}$  as a regular function in  $z \in \mathbb{C} \setminus [e_{\min}, e_{\max}]$  is defined as

$$\Delta(\lambda, \mu; z) = (1 - \lambda b(z))^2 - \lambda^2 a^2(z) - \mu a(z).$$

$$c(z) = a^2(z) - b^2(z).$$

For any  $z \in \mathbb{R} \setminus [e_{\min}, e_{\max}]$ , the number  $a(z)$  is non-zero, and the equality

$$\Delta(\lambda, \mu; z) = a(z) \left( \frac{1}{a(z)} - \frac{2b(z)}{a(z)} \lambda - \frac{c(z)}{a(z)} \lambda^2 - \mu \right)$$

allows us to consider the parabola

$$P_z(\lambda, \mu) := \frac{1}{a(z)} - \frac{2b(z)}{a(z)} \lambda - \frac{c(z)}{a(z)} \lambda^2 - \mu = 0$$

in the  $\lambda\mu$ -plane with  $\lambda$ -intercepts

$$\kappa(z) = \frac{1}{b(z) + a(z)} \quad \text{and} \quad \nu(z) = \frac{1}{b(z) - a(z)}$$

instead of the equality  $\Delta(\lambda, \mu; z) = 0$ .

**Lemma 1.** The number  $z \in \mathbb{C} \setminus [e_{\min}, e_{\max}]$  is an eigenvalue of  $H_{\lambda\mu}$  if and only if  $\Delta(\lambda, \mu; z) = 0$ .

**Lemma 2.** The parabola  $P_z(\lambda, \mu) = 0$  can be defined at  $e_{\max}$  as

$$P_{e_{\max}}(\lambda, \mu) := -2\lambda + \frac{2}{\kappa_r} \lambda^2 - \mu = 0 \quad \text{for } d=1, 2 \text{ with } |x_0|_1 - \text{even},$$

$$P_{e_{\max}}(\lambda, \mu) := 2\lambda - \frac{2}{\nu_r} \lambda^2 - \mu = 0 \quad \text{for } d=1, 2 \text{ with } |x_0|_1 - \text{odd},$$

$$P_{e_{\max}}(\lambda, \mu) := \frac{1}{a_r} - \frac{2b_r}{a_r} \lambda - \frac{c_r}{a_r} \lambda^2 - \mu = 0 \quad \text{for } d \geq 3,$$

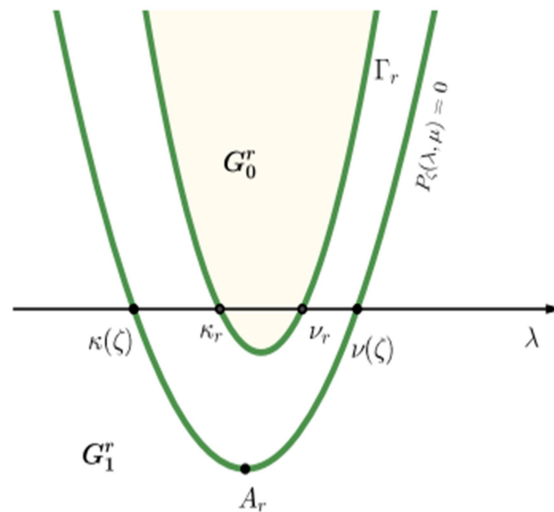
where

$$a_r := \lim_{z \rightarrow e_{\max}^+} a(z), \quad b_r := \lim_{z \rightarrow e_{\max}^+} b(z), \quad c_r := a_r^2 - b_r^2.$$

$$\kappa_r = \lim_{z \rightarrow 2d^+} \kappa_r(z), \quad \nu_r = \lim_{z \rightarrow 2d^+} \nu_r(z)$$

Let  $\Gamma_r$  be graph of the parabola  $P_{e_{\max}}(\lambda, \mu) = 0$  in the  $\lambda\mu$ -plane. It divides the  $(\lambda, \mu)$  plane into two disjoint parts

$$G_0^r = \{(\lambda, \mu) : P_{e_{\max}}(\lambda, \mu) > 0\}, \quad G_1^r = \{(\lambda, \mu) : P_{e_{\max}}(\lambda, \mu) < 0\}.$$



**Figure :** Location of  $P_{\zeta}(\lambda, \mu) = 0$ ,  $\zeta > e_{\max}$ .

**Theorem 1.**

(a) For  $(\lambda, \mu) \in G_0^r \cup \Gamma_r$ , the operator  $H_{\lambda\mu}$  has no eigenvalues in  $[e_{\max}, \infty)$ .

(b) For  $(\lambda, \mu) \in G_1^r$ , the operator  $H_{\lambda\mu}$  has a simple eigenvalue in  $[e_{\max}, \infty)$ .

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**CHEBYSHEV KO‘PHADLARI ORQALI FUNKSIYALARNI YAQNLASHTIRISH**

**Jo’liiyeva Dildora G’ayrat qizi**

Denov tadbirkorlik va pedagogika instituti  
1-kurs magistranti tel.: 998910778103